

Физические механизмы спин-
токовой модели поляризации
мультиферроиков

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Moscow, Russia - 2024

Multiferroic materials
-- Ferromagnetic materials with
Ferroelectric properties

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graph TD; A["Multiferroic materials  
-- Ferromagnetic materials with  
Ferroelectric properties"] --> B["I-st Type  
Coexistence of  
Ferromagnetic  
and  
Ferroelectric  
properties"]; A --> C["II-nd Type  
Multiferroics  
of  
Spin origin"];
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I-st Type
Coexistence of
Ferromagnetic
and
Ferroelectric
properties

II-nd Type
Multiferroics
of
Spin origin

Multiferroic materials, II-nd Type -- Classification

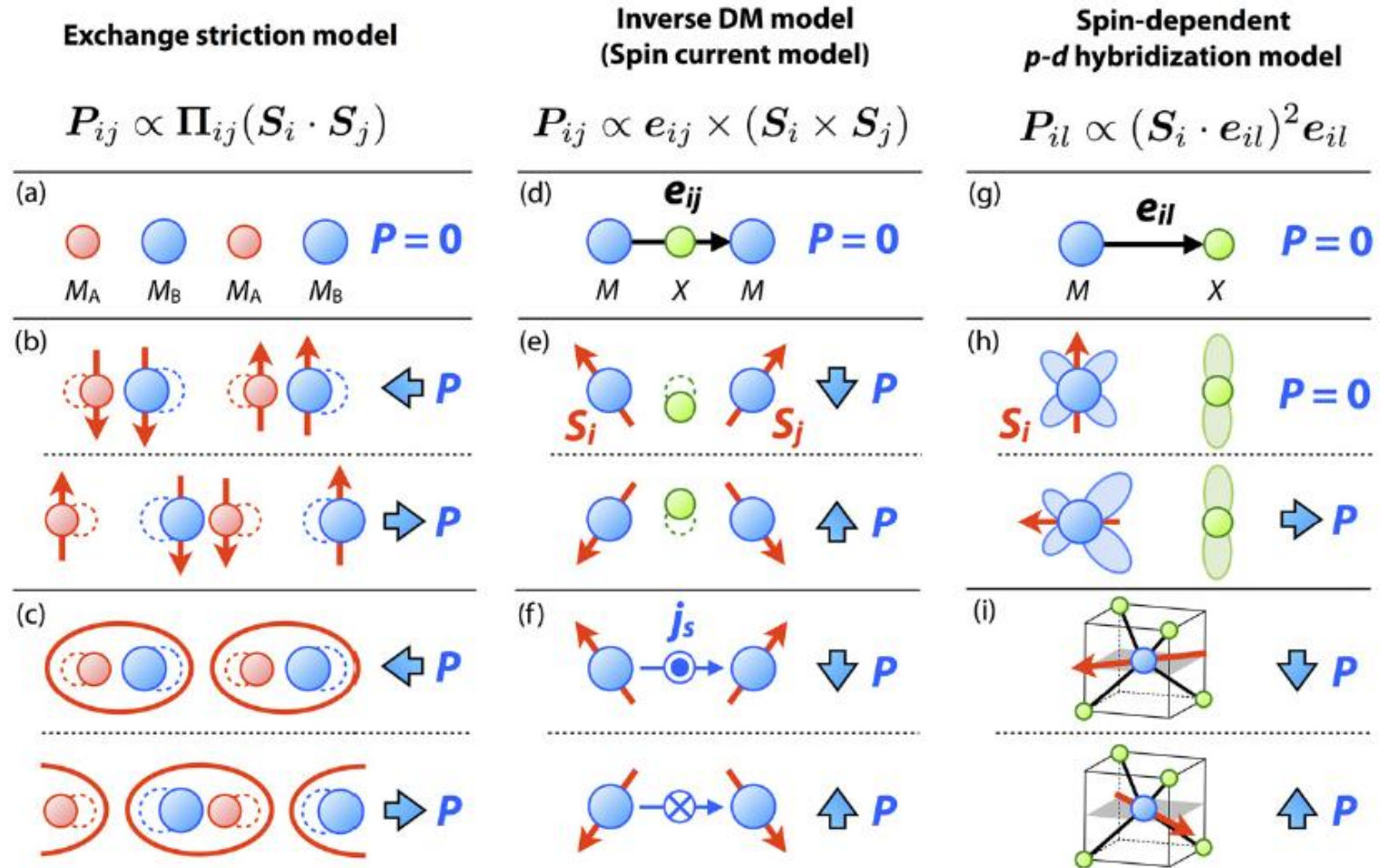


Figure 2. Three major mechanisms of ferroelectricity of spin origin: (a)–(c) exchange-striction mechanism arising from the symmetric spin

Models of Multiferroic materials

Exchange striction model

$$P_{ij} \propto \Pi_{ij} (\mathbf{S}_i \cdot \mathbf{S}_j)$$

**Inverse DM model
(Spin current model)**

$$P_{ij} \propto \mathbf{e}_{ij} \times (\mathbf{S}_i \times \mathbf{S}_j)$$

**Spin-dependent
p-d hybridization model**

$$P_{il} \propto (\mathbf{S}_i \cdot \mathbf{e}_{il})^2 \mathbf{e}_{il}$$

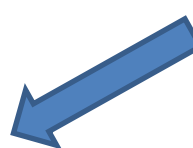
$$P^\alpha = g_{0\Pi}^\alpha \mathbf{S}^2 + \frac{1}{6} g_{\Pi}^\alpha (\mathbf{S} \cdot \Delta \mathbf{S}),$$

Spin current model:

$$P^\beta \sim \varepsilon^{\beta\mu\nu} J^{\mu\nu}$$

J is the spin current tensor

Tokura et al , **Multiferroics of spin origin**, Rep. Prog. Phys. **77**, 076501 (2014)
 S. Dong et al , Multiferroic materials and magnetoelectric physics:
 symmetry, entanglement, excitation, and topology, Advances in Physics, 2015



$$P^\beta = \alpha \left(M^\beta (\nabla \cdot \mathbf{M}) - (\mathbf{M} \cdot \nabla) M^\beta \right)$$

Mostovoy M., **Ferroelectricity in Spiral Magnets**, PRL **96**, 067601 (2006)

Model and corresponding spin evolution equation I

$$\partial_t \mathbf{S} = \frac{2\mu}{\hbar} [\mathbf{S} \times \mathbf{B}] + \frac{1}{6} g_u [\mathbf{S} \times \Delta \mathbf{S}] + \frac{2\mu}{\hbar c} \varepsilon^{\alpha\beta\gamma} \varepsilon^{\beta\mu\nu} J^{\gamma\mu} E^\nu$$

$$+ \frac{1}{3} g_{(\beta)} [(\mathbf{S} \cdot [\boldsymbol{\delta} \times \nabla]) \mathbf{S} - S^\beta [\boldsymbol{\delta} \times \nabla] S^\beta]$$

$$J^{\gamma\mu} \sim S^\gamma \cdot v^\mu \quad \mathbf{S} \equiv S^\alpha$$

$$\mathbf{S}(\mathbf{r}, t) = \int \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \Psi^\dagger(R, t) \hat{\mathbf{s}}_i \Psi(R, t) d^{3N} R$$

P. A. Andreev, M. I. Trukhanova, arXiv:2312.16321

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MAGNETIC SOLITONS

A.M. KOSEVICH, B.A. IVANOV and A.S. KOVALEV

average spin of a lattice site,

$$M_n = -\frac{2\mu_0}{a^3} \langle S_n \rangle = -\frac{2\mu_0}{a^3} S(x_n), \quad S_n \rightarrow -\frac{a^3}{2\mu_0} M(x_n),$$

where a is the interaction distance.

If the angle between adjacent spin vectors is small, then the vectors as continuous functions of the x coordinate, and we can write the fol

$$M_{n+n_0} = M(x_n) + x_i(n_0) \frac{\partial M}{\partial x_i} + \frac{1}{2} x_i(n_0) x_k(n_0) \frac{\partial^2 M}{\partial x_i \partial x_k} + \dots$$

Using the procedure prescribed by eqs. (1.9) and (1.10) for the real characteristics of a discrete spin system to a continuous description of the expression for the energy density of a ferromagnet corresponding writing the energy density, it should be remembered that the exchange and that the values of J_1 , J_2 , and J_3 are quite close to each other. exchange constants J_α , we may assume that

$$J_\alpha = J_0 + a^2 j_\alpha, \quad \alpha = 1, 2, 3.$$

To order a^2 , one can write

$$w = \frac{\alpha}{2} \frac{\partial M}{\partial x_k} \cdot \frac{\partial M}{\partial x_k} + w_{\text{an}} - M \cdot H, \quad \alpha = J_0 a^5 / (2\mu_0)^2,$$

$$\mathbf{S}(\mathbf{r}, t) = \int d^{3N} R \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \cdot$$

$$\Psi^\dagger(R, t) \hat{\mathbf{s}}_i \Psi(R, t)$$

$$\frac{\partial}{\partial t} \mathbf{S}(\mathbf{r}, t) = \frac{i}{\hbar} \int d^{3N} R \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \cdot$$

$$\Psi^\dagger(R, t) [\hat{H}, \hat{\mathbf{s}}_i] \Psi(R, t)$$

Model and corresponding spin evolution equation II

$$\partial_t \mathbf{S} = \frac{2\mu}{\hbar} [\mathbf{S} \times \mathbf{B}] + \frac{1}{6} g_u [\mathbf{S} \times \Delta \mathbf{S}] + \frac{2\mu}{\hbar c} \varepsilon^{\alpha\beta\gamma} \varepsilon^{\beta\mu\nu} J^{\gamma\mu} E^\nu$$

$$+ \frac{1}{3} g_{(\beta)} \left[(\mathbf{S} \cdot [\boldsymbol{\delta} \times \nabla]) \mathbf{S} - S^\beta [\boldsymbol{\delta} \times \nabla] S^\beta \right]$$

$$J^{\gamma\mu} \sim S^\gamma \cdot v^\mu \quad \mathbf{S} \equiv S^\alpha$$

$$g_u = \int \xi^2 U(\xi) d^3 \xi \quad g_{(\beta)} = \int \xi^2 \beta(\xi) d^3 \xi \quad \mathbf{D}_{ij} = \beta(r_{ij}) \cdot \mathbf{r}_{ij} \times \boldsymbol{\delta}$$

$$\hat{H} = \sum_i \left[-\hat{\mathbf{d}}_i \cdot \mathbf{E}_i - \hat{\boldsymbol{\mu}}_i \cdot \mathbf{B}_i - \frac{1}{m_i c} (\hat{\boldsymbol{\mu}}_i \cdot [\mathbf{E}_i \times \hat{\mathbf{p}}_i]) \right.$$

$$\left. - \frac{1}{2} \sum_{j \neq i} (U_{ij} \hat{\mathbf{s}}_i \cdot \hat{\mathbf{s}}_j + \mathbf{D}_{ij} \cdot [\hat{\mathbf{s}}_i \times \hat{\mathbf{s}}_j]) \right]$$

$$i\hbar \partial_t \Psi(R, t) = \hat{H} \Psi(R, t)$$

Same model and corresponding Momentum evolution equation –derivation of the Spin Current model

$$\hat{H} = \sum_i \left[-\hat{\mathbf{d}}_i \cdot \mathbf{E}_i - \hat{\boldsymbol{\mu}}_i \cdot \mathbf{B}_i - \frac{1}{m_i c} (\hat{\boldsymbol{\mu}}_i \cdot [\mathbf{E}_i \times \hat{\mathbf{p}}_i]) \right. \\ \left. - \frac{1}{2} \sum_{j \neq i} (U_{ij} \hat{\mathbf{s}}_i \cdot \hat{\mathbf{s}}_j + D_{ij} \cdot [\hat{\mathbf{s}}_i \times \hat{\mathbf{s}}_j]) \right] \quad i\hbar \partial_t \Psi(R, t) = \hat{H} \Psi(R, t)$$

$$\partial_t \mathbf{p} = \partial_t (m n \mathbf{v}) = P^\beta (\nabla E^\beta) + \mu S^\beta \nabla B^\beta + \frac{\mu}{m c} \varepsilon^{\beta\gamma\delta} J^{\delta\gamma} (\nabla E^\beta) \\ + g_{0u} S^\beta \nabla S^\beta + \frac{1}{3} g_{(\beta)} [(\boldsymbol{\delta} \cdot \mathbf{S}) \nabla (\nabla \cdot \mathbf{S}) - (\mathbf{S} \cdot \nabla) \nabla (\boldsymbol{\delta} \cdot \mathbf{S})]$$

$$g_{0u} = \int U(\xi) d^3 \xi$$

$$g_{(\beta)} = \int \xi^2 \beta(\xi) d^3 \xi$$

$$\mathbf{D}_{ij} = \beta(r_{ij}) \mathbf{r}_{ij} \times \boldsymbol{\delta}$$

Spin Current model

As a condition of absence of currents

$$P^\beta = \frac{\mu}{mc} \varepsilon^{\beta\gamma\delta} J^{\gamma\delta}$$

Spin Current: How to get it?

$$\partial_t \mathbf{S} = \frac{2\mu}{\hbar} [\mathbf{S} \times \mathbf{B}] + \frac{1}{6} g_u [\mathbf{S} \times \Delta \mathbf{S}] + \frac{2\mu}{\hbar c} \varepsilon^{\alpha\beta\gamma} \varepsilon^{\beta\mu\nu} J^{\gamma\mu} E^\nu + \frac{1}{3} g_{(\beta)} [(\mathbf{S} \cdot [\delta \times \nabla]) \mathbf{S} - S^\beta [\delta \times \nabla] S^\beta]$$

$$g_u = \int \xi^2 U(\xi) d^3 \xi$$

$$g_{(\beta)} = \int \xi^2 \beta(\xi) d^3 \xi$$

$$\mathbf{D}_{ij} = \beta(r_{ij}) \mathbf{r}_{ij} \times \delta$$

$$\frac{1}{6} g_u \varepsilon^{\alpha\beta\gamma} S^\beta \Delta S^\gamma = -\partial^\beta J_{HH}^{\alpha\beta}$$

$$J_{HH}^{\alpha\beta} = -\frac{1}{6} g_u \varepsilon^{\alpha\mu\nu} S^\mu \partial^\beta S^\nu$$

Polarization for the Spin Current from HH

$$P^\beta = \frac{\mu}{mc} \varepsilon^{\beta\gamma\delta} J^{\gamma\delta}$$

$$\frac{1}{6} g_u \varepsilon^{\alpha\beta\gamma} S^\beta \Delta S^\gamma = -\partial^\beta J_{HH}^{\alpha\beta}$$

$$J_{HH}^{\alpha\beta} = -\frac{1}{6} g_u \varepsilon^{\alpha\mu\nu} S^\mu \partial^\beta S^\nu$$

$$\begin{aligned} P_{HH}^\beta &= \frac{\mu}{mc} \varepsilon^{\beta\gamma\delta} J_{HH}^{\gamma\delta} \\ &= \frac{1}{6} \frac{\mu}{mc} g_u \left((\mathbf{S} \cdot \nabla) S^\beta - S^\beta (\nabla \cdot \mathbf{S}) \right) \end{aligned}$$

Same polarization can be found
as average of edm operator

$$\mathbf{P}(\mathbf{r}, t) = \int \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \Psi^\dagger(R, t) \cdot \hat{\mathbf{d}}_i \cdot \Psi(R, t) d^{3N} R$$

$$\hat{\mathbf{d}}_i = \sum_{j \neq i} \alpha_{ij}(r_{ij}) [\mathbf{r}_{ij} \times [\hat{\mathbf{s}}_i \times \hat{\mathbf{s}}_j]] \quad \alpha_{ij} = \frac{\mu}{2mc} U_{ij} \quad U_{ij} \sim \text{exchange int}$$

Polarization for the Spin Current from HH:

Comparison with “Tokura and Mostovoy”

$$P^{\beta}_{HH} = \frac{\mu}{mc} \varepsilon^{\beta\gamma\delta} J^{\gamma\delta}_{HH} = \frac{1}{6} \frac{\mu}{mc} g_u \left((\mathbf{S} \cdot \nabla) S^{\beta} - S^{\beta} (\nabla \cdot \mathbf{S}) \right)$$

Same polarization can be found
as average of edm operator

$$\mathbf{P}(\mathbf{r}, t) = \int \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \Psi^{\dagger}(R, t) \cdot \hat{\mathbf{d}}_i \cdot \Psi(R, t) d^{3N} R$$

$$\hat{\mathbf{d}}_i = \sum_{j \neq i} \alpha_{ij}(r_{ij}) [\mathbf{r}_{ij} \times [\hat{\mathbf{s}}_i \times \hat{\mathbf{s}}_j]] \quad \alpha_{ij} = \frac{\mu}{2mc} U_{ij} \quad U_{ij} \sim \text{exchange int}$$

Spin Current from DMI

$$P^\beta = \frac{\mu}{mc} \varepsilon^{\beta\gamma\delta} J^{\gamma\delta}$$

$$\partial_t \mathbf{S} = \frac{2\mu}{\hbar} [\mathbf{S} \times \mathbf{B}] + \frac{1}{6} g_u [\mathbf{S} \times \Delta \mathbf{S}] + \frac{2\mu}{\hbar c} \varepsilon^{\alpha\beta\gamma} \varepsilon^{\beta\mu\nu} J^{\gamma\mu} E^\nu$$

$$+ \frac{1}{3} g_{(\beta)} \left[(\mathbf{S} \cdot [\boldsymbol{\delta} \times \nabla]) \mathbf{S} - S^\beta [\boldsymbol{\delta} \times \nabla] S^\beta \right]$$

$$g_u = \int \xi^2 U(\xi) d^3 \xi$$

$$g_{(\beta)} = \int \xi^2 \beta(\xi) d^3 \xi$$

$$\mathbf{D}_{ij} = \beta(r_{ij}) \mathbf{r}_{ij} \times \boldsymbol{\delta}$$

$$-\frac{1}{3} g_{(\beta)} S^\mu [\boldsymbol{\delta} \times \nabla] S^\mu \equiv -\partial^\beta J_{DM}^{\alpha\beta}$$

$$J_{DM}^{\alpha\beta} = -\frac{1}{6} g_{(\beta)} \varepsilon^{\alpha\beta\gamma} \delta^\gamma (S^\mu S^\mu)$$

$$P_{DM}^\beta = \frac{\mu}{mc} \varepsilon^{\beta\gamma\delta} J_{DM}^{\gamma\delta}$$

$$= -\frac{1}{3} \frac{\mu}{mc} g_{(\beta)} \delta^\beta \cdot \mathbf{S}^2$$

Spin Current from DMI: Comparison with “Tokura”

$$P^\beta = \frac{\mu}{mc} \varepsilon^{\beta\gamma\delta} J^{\gamma\delta}$$

$$J_{DM}^{\alpha\beta} = -\frac{1}{6} g_{(\beta)} \varepsilon^{\alpha\beta\gamma} \delta^\gamma \cdot \mathbf{S}^2$$

$$\begin{aligned} P_{DM}^\beta &= \frac{\mu}{mc} \varepsilon^{\beta\gamma\delta} J_{DM}^{\gamma\delta} \\ &= -\frac{1}{3} \frac{\mu}{mc} g_{(\beta)} \delta^\beta \cdot \mathbf{S}^2 \end{aligned}$$

Same polarization can be found
as average of edm operator

$$\mathbf{P}(\mathbf{r}, t) = \int \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \Psi^\dagger(R, t) \cdot \hat{\mathbf{d}}_i \cdot \Psi(R, t) d^{3N} R$$

$$\hat{\mathbf{d}}_i = \sum_{j \neq i} \mathbf{\Pi}_{ij}(r_{ij}) (\hat{\mathbf{s}}_i \cdot \hat{\mathbf{s}}_j)$$

$$\mathbf{\Pi}_{ij} = -\frac{\mu}{3mc} r_{ij}^2 \beta_{ij} \boldsymbol{\delta}$$

$$\frac{\partial \mathbf{\Pi}_{ij}}{\partial r_{ij}} = \frac{\mu}{mc} r_{ij} \beta_{ij} \boldsymbol{\delta}$$

Spin Current from DMI: non-homogeneous contribution

$$\begin{aligned}
 P_{DM}^\beta &= \frac{\mu}{mc} \varepsilon^{\beta\gamma\delta} J_{DM}^{\gamma\delta} \\
 &= -\frac{1}{3} \frac{\mu}{mc} \delta^{\beta\gamma} \left(g_{(\beta)} \mathbf{S}^2 + \frac{1}{5} g_{2(\beta)} \mathbf{S} \cdot \Delta \mathbf{S} \right)
 \end{aligned}$$

Same polarization can be found
as average of edm operator

$$\mathbf{P}(\mathbf{r}, t) = \int \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \Psi^\dagger(R, t) \cdot \hat{\mathbf{d}}_i \cdot \Psi(R, t) d^{3N} R$$

$$\hat{\mathbf{d}}_i = \sum_{j \neq i} \mathbf{\Pi}_{ij}(r_{ij}) (\hat{\mathbf{s}}_i \cdot \hat{\mathbf{s}}_j)$$

$$\mathbf{\Pi}_{ij} = -\frac{\mu}{3mc} r_{ij}^2 \beta_{ij} \boldsymbol{\delta}$$

$$\frac{\partial \mathbf{\Pi}_{ij}}{\partial r_{ij}} = \frac{\mu}{mc} r_{ij} \beta_{ij} \boldsymbol{\delta}$$

Spin evolution equation and Polarization evolution equation

$$\mathbf{S}(\mathbf{r}, t) = \int \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \Psi^\dagger(R, t) \hat{\mathbf{s}}_i \Psi(R, t) d^{3N} R$$

$$\mathbf{P}(\mathbf{r}, t) = \int \sum_i \delta(\mathbf{r} - \mathbf{r}_i) \Psi^\dagger(R, t) \cdot \hat{\mathbf{d}}_i \cdot \Psi(R, t) d^{3N} R$$

$$\hat{\mathbf{d}}_i = \sum_{j \neq i} \alpha_{ij}(r_{ij}) [\mathbf{r}_{ij} \times [\hat{\mathbf{s}}_i \times \hat{\mathbf{s}}_j]]$$

$$\hat{H} = \sum_i (-\hat{\boldsymbol{\mu}}_i \cdot \mathbf{B}_i)$$

$$\partial_t P^\alpha = \frac{1}{3} \varepsilon^{\alpha\beta\gamma} \gamma g_{(\alpha)} \cdot \left[S^2 \cdot \partial^\beta B^\gamma + S^\gamma S^\mu \cdot \partial^\beta B^\mu + B^\mu (S^\gamma \partial^\beta S^\mu - S^\mu \partial^\beta S^\gamma) \right]$$

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Thank you!